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A straightforward kinematic analysis of diffraction from metastable three-dimensional crystallites of Ge grown on Si(001) is presented. Low-energy electron diffraction data from these crystallites are shown to agree with diffraction images calculated for a structure determined from scanning-tunneling microscopy data. Additionally, reflection high-energy electron diffraction images predicted for these crystals agree with existing data.

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Diffraction Determination of the Structure of Metastable  
Three-Dimensional Crystals of Ge Grown On Si(001)

by

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## Introduction

Recent scanning tunneling microscopy (STM) measurements of the growth of Ge on Si(001) have revealed metastable Ge crystallites that form on top of the initial  $\sim$  three-monolayer two-dimensional Ge phase.<sup>(1)</sup> Figure 1 is an STM scan of one of these metastable "hut" clusters, showing its well-defined facet structure. Each crystallite consists of a pyramid or elongated pyramid with four equivalent facets. The facets were proposed<sup>(1)</sup>, on the basis of detailed STM scans, to be reconstructed {105} planes consisting of (001) terraces separated by monatomic-height steps running parallel to  $\langle 100 \rangle$  directions. Each terrace is two dimer pairs wide with the dimer pairs angled  $45^\circ$  with respect to the step and rotating  $90^\circ$  upon crossing a step. Alternatively, a facet can be described as a vicinal (001) surface having a miscut of  $11.3^\circ$  toward  $\langle 100 \rangle$ .

We have also observed this phase using low-energy electron diffraction (LEED). Diffraction patterns of Si(001) with and without the Ge hut crystals present are shown in Fig. 2. The Ge huts add extra features about the substrate beams that look like X's. These diffraction patterns are relatively straightforward to interpret in terms of the structure causing them. Vicinal surfaces are well known to cause a splitting of diffracted beams in a direction normal to the steps<sup>(2)</sup>. Because the extra spots in Fig. 2b are split in a direction  $45^\circ$  relative to the symmetry of the substrate reflections (i.e., the substrate unit mesh has sides parallel to

<110>) it is immediately known that the facets must have <100>-type zone axes, i.e., steps parallel to <100> as in the model established by STM.<sup>1</sup> The diffraction pattern in Fig. 2b can be reproduced via a kinematic calculation for the unit mesh proposed in Ref. 1. Taking the dimer pairs from odd-and even-level terraces to scatter equivalently yields the structure factor:

$$F = \cos(\vec{S} \cdot (\hat{j} - \hat{i})a/4) + \cos(\vec{S} \cdot (\hat{i} + \hat{j})a/4)e^{i\vec{S} \cdot \vec{c}}, \quad (1)$$

where the basis vector  $\vec{c} = 5\hat{a}_i/4 + 3\hat{a}_j/4 - \hat{a}_k/4$ ,  $a = 5.66\text{\AA}$  is the bulk lattice constant for Ge, and  $\vec{S}$  is a reciprocal-lattice vector. The scattered intensity from this (105) facet is then:

$$I = \left| \sum_{n=0}^N e^{in\vec{S} \cdot \vec{a}} \sum_{m=0}^M e^{im\vec{S} \cdot \vec{b}} \right|^2 \quad (2)$$

with unit mesh vectors  $\vec{a} = 5\hat{a}_i/2 - \hat{a}_k/2$ ,  $\vec{b} = 2\hat{a}_j$  and  $n, m$  integers. For simplicity the facets are assumed to be rectangular with length  $M|\vec{b}|$  and width  $N|\vec{a}|$ . The scattered intensity from a cluster is taken to be an incoherent sum of the intensities from each of its four facets. Figure 3a shows the resulting diffraction pattern at the diffraction conditions used to obtain the pattern shown in Fig. 2b. It can be seen that there is a 1:1 match between the observed and calculated diffraction patterns. The magnitude of the facet splitting relative to the substrate beams is set completely by the diffraction conditions, while the size and shape of the facet

beams are determined by the distribution of cluster sizes and shapes. Different facet structures give completely different patterns. For example, the diffraction pattern from a cluster with (117) facets is shown in Fig. 3b.

Figure 4 shows diffraction patterns for Ge clusters with (105) facets at diffraction conditions appropriate to reflection high-energy electron diffraction (RHEED). Patterns in two azimuthal orientations are shown. These two patterns illustrate the importance, in experiments, of changing the azimuthal direction of the incident-electron beam with respect to the substrate. In this example, changing the azimuth produces highly different patterns, from which the facet orientation can be identified. RHEED patterns from facets are most easily interpreted when the electron beam is directed along the steps or zone axes of the facet, as in Fig. 4a. Under such conditions diffraction from the facet 'rods' yields lines that clearly connect 3D diffraction features. The angle between these lines and the shadow edge defines the facet angle, the angle between the substrate surface and the facet. For electron-beam azimuths away from a facet's zone axis the projected angle between the facet 'rods' and shadow edge is less than the facet angle because of the way in which the Ewald sphere cuts these 'rods'. For example, it would be incorrectly concluded that the facet angle is  $8.0^\circ$  instead of  $11.3^\circ$  if measurements of (105) facets were made with the electron beam along [110] instead of along [100].

The metastable "hut" crystals appear to have been observed by others using RHEED<sup>3</sup>. Their measurements were, however, interpreted in terms of diffraction from 3D clusters having a structure suggested by Ourmazd and Bean.<sup>4</sup> This structure, consisting of alternating Si(111) and Ge(111) layers, yields a facet with a  $\langle 110 \rangle$  zone axis. Such facets produce diffraction patterns that have no resemblance to these shown in Fig. 4; the experimental patterns<sup>3</sup> do, on the other hand, look like those shown in Fig. 4, suggesting that the phase observed by the authors of Ref. 3 is the same as ours. Although our calculations do not include multiple scattering, these effects cannot change the diffraction pattern in its essential features.

To our knowledge, this is the first time that a structure initially proposed on the basis of STM was subsequently confirmed by diffraction, rather than vice versa.

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### Figure Captions

- Fig. 1 STM image of a metastable Ge crystal. These crystals form on Si(001) for coverages  $>3$  monolayers and for  $T < 800\text{K}$ . The crystallite is  $\sim 30\text{\AA}$  high and  $\sim 200\text{\AA}$  long. From ref. 1.
- Fig. 2 LEED patterns of clean and Ge-covered Si(001). a) LEED pattern of clean Si(001) showing a two-domain  $(2\times 1)$  diffraction pattern. Diffracted beams lie along  $[110]$ . b) LEED pattern showing the extra reflections from Ge crystallite facets, obtained with a  $\sim 100\text{eV}$  beam at an angle of incidence of  $\sim 80^\circ$  measured from the horizontal, with the beam directed along  $[110]$ .
- Fig. 3 Simulated LEED pattern for clusters with two different facet orientations: a)  $\{105\}$  facets. b)  $\{117\}$  facets.  $\{117\}$  facets were used because they have a facet angle of  $11.4^\circ$ , which is very close to that of the  $\{105\}$  facets, but in a different orientation. The facet features for  $\{117\}$  facets are split in directions that are along lines connecting substrate beams. Both pictures are calculated at the diffraction conditions used in Fig. 2, assuming all clusters have the same size and  $(N,M) = (7,12)$  in Eq. (2).
- Fig. 4 Simulated 20 KeV RHEED patterns for diffraction conditions satisfying the  $(004)$  reflection, with the beam directed along a)  $[100]$  and b)  $[110]$  for clusters with  $\{105\}$  facets. The  $(004)$  reflection at this energy occurs at an angle of incidence of  $1.8^\circ$ .









